

Chapter 12B

HEAT TRANSFER AND FLOW PHENOMENA II

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12B.1 Introduction

In forced convective boiling the boiling crisis* occurs when the heat flux is raised to such a high level that the heated surface can no longer support continuous liquid contact. This heat flux is usually referred to as the critical heat flux (CHF) or dryout heat flux. It is characterized either by a sudden rise in surface temperature caused by blanketing of the heated surface by a stable vapour layer, or by small surface temperature spikes corresponding to the appearance and disappearance of dry patches.

Failure of the heated surface may occur once the CHF is exceeded. High surface temperatures can be avoided by either operating at heat flux levels well below the critical heat flux or by operating at conditions where the post-dryout heat transfer is reasonably effective in keeping the surface temperatures at moderate levels. It is common practice in the thermal design of reactors to avoid dryout.

12B.2 Critical Heat Flux

12B.2.1 General

This chapter discussed the critical heat flux (CHF) for pool boiling or axial flow inside tubes or rod bundles. Axial flow CHF has been extensively studied and many excellent reviews are available (e.g., Tong, 1972; Hewitt, 1978; and Collier, 1980). Much effort has been spent on correlating the available CHF data; it is conservatively estimated that over 400 CHF correlations are currently available. This proliferation of CHF correlations illustrates the sad state-of-the-art in modelling the CHF phenomenon.

12B.2.2 Dryout Mechanisms

In flow boiling the dryout mechanisms may depend strongly on flow regime and phase distributions which in turn are controlled by pressure, flow rate and quality. Figure 1 shows the variation in CHF and dryout mechanisms with quality. The following dryout mechanisms are expected to be the most important ones:

Nucleation induced This type of dryout is encountered at high subcooling where heat is transferred very efficiently by nucleate boiling; bubbles grow and collapse at the wall. Some convection will take place between the bubbles.

* Other terms used to denote the boiling crisis: burnout, dryout, departure from nucleate boiling (DNB).

Nucleation induced dryout occurs at very high surface heat fluxes. It has been suggested (Collier, 1980; Tong, 1972) that dryout is due to the spreading of drypath following microlayer evaporation under a bubble and coalescence of adjacent bubbles although no definite proof of this is yet available. The occurrence of dryout here only depends on the local surface heat flux and flow conditions and is not affected by upstream heat flux distribution. The surface temperature rise under these conditions is very rapid.

Bubble clouding In subcooled and saturated nucleate boiling the number of bubbles generated depends on the heat flux and bulk temperature. The bubble population density near the heated surface increases with increasing heat flux and a so-called bubble boundary layer (Tong, 1965, 1972) often forms a short distance away from the surface. If this layer is sufficiently thick it can impede the flow of coolant to the heated surface. This in turn leads to a further increase in bubble population until the wall becomes so hot that a vapour patch forms over the heated surface. This type of boiling crisis is also characterized by a fast rise of the heated surface temperature.

Annular Film Dryout In the annular dispersed flow regime (high void fraction flow) the liquid will be in the form of a liquid film covering the walls and entrained droplets moving at a higher velocity in the core. Continuous thinning of the liquid film will take place due to the combined effect of entrainment and evaporation. At the dryout location the liquid film eventually breaks down. The temperature rise accompanying this film breakdown is usually moderate.

12B.2.3 Prediction Methods

CHF prediction methods may be subdivided into three types, i.e.:

(i) Local conditions type CHF correlations. $CHF = f(P, G, X \text{ or } \alpha, \text{ cross section geometry})$. These correlations are convenient to use for predicting location and magnitude of the CHF. Effects of axial flux distribution, spacers, flux spikes and flow transients (e.g., flow stagnation) usually require a local conditions approach, sometimes combined with techniques which consider the upstream flow history (e.g., boiling length average approach, F-factor method, grid spacer factor).

(ii) Global conditions type correlations. The burnout power = $f(P, G, H_{in}, L_h, \text{ cross-section geometry})$. These correlations predict only the burnout power; they cannot be used to predict the location or magnitude of the CHF. Also, these correlations are incapable of considering the effect of a variation in axial flux distribution. They are primarily used to predict burnout power for a given geometry and axial flux distribution.

(iii) Standard CHF table method. As most empirical correlations and analytical models have a limited range of application, the need for a more general technique is obvious. Attempts have been made in the USSR to construct a standard table of CHF values for a given geometry (Doroshchuk, 1975). The table approach has been continued at Chalk River, at CEN-Grenoble and at the University of Ottawa using the Chalk River data base (10,000 tube CHF data). Table 1 shows the CHF table for an 8 mm tube and discrete values for P, G and X. Note that this table has the correct asymptotic trends (e.g., $CHF = 0$ at $X = 1.0$ and $CHF = CHF_{PB}$ at $G = 0$).

12B.2.4 Discussion

Most CHF correlations are based on tube data. In general, the CHF in tubes is greater than the corresponding value in bundles because:

- (a) the subchannel with the highest enthalpy usually reaches the boiling crisis before a subchannel at cross section average conditions
- (b) of the cold wall effect: a cold wall reduces liquid available for cooling heated rods (Tong, 1972b)
- (c) inside a subchannel, significant variation in velocities and near wall void fraction may occur (Groeneveld, 1973). This is particularly evident in the gap between adjacent rods of tightly spaced rod bundles. One can thus expect preferential CHF locations within a single subchannel
- (d) upstream or fluid memory effects can become very important, especially in the net quality region. Large differences in subchannel enthalpy rise rates can result in significant asymmetric CHF patterns (McPherson, 1971).

Some of the above effects can be partially incorporated in subchannel codes which predict the flow, enthalpy and pressure in each subchannel. An excellent review of subchannel type analyses has been presented by Weissman (1975). Caution must be exercised when using subchannel codes as, frequently, verification data are lacking. Subchannel CHF correlations, based on bundle CHF data and predicted subchannel conditions are especially unreliable. Their application should be limited to the narrow range of their data base. Sub-channel codes currently under development are based on 2- or 3-fluid models, (e.g., Tahir, 1983); they are expected to be capable of predicting the bundle CHF with a much better accuracy than their predecessors.

Tube and annuli CHF correlations may be useful in predicting the CHF behaviour of bundles, especially for well-balanced bundles having wide rod spacings. Correlation factors should then be used to include specific bundle characteristics (Groeneveld, 1974).

12B.3 Transition Boiling

12B.3.1 General

Transition boiling is a rather unique heat transfer mode because, here, the heat flux generally decreases with an increase in surface temperature. It is basically a combination of unstable film boiling and unstable nucleate boiling alternately existing at any given location on a heating surface.

The transition boiling section of the boiling curve is bounded by the critical heat flux and the minimum heat flux (Fig. 2). The critical heat flux has been extensively studied but, there are still wide ranges of conditions where data are virtually nonexistent. The minimum heat flux has undergone less study; it is known to be affected by flow conditions, fluid properties and heated surface properties.

In pool boiling, at surface temperatures just above the boiling crisis temperature, the heated surface is partially covered with unstable vapor patches, varying with space and time. The formation of such dry patches is accompanied by a drastic reduction in heat transfer coefficient corresponding to the change from nucleate boiling to film boiling; the corresponding reduction in local vapor generation permits the liquid to momentarily rewet the heated surface.

(kW/m^2)

12B-4a

TABLE 1 : STANDARD CHF TABLE (cont'd)

PRESSURE (KPA)	G (KG/M ² /S)	QUALITY																
		-0.15	-0.10	-0.05	0.00	0.05	0.10	0.15	0.20	0.25	0.30	0.40	0.50	0.60	0.70	0.80	0.90	0.0
1500	0	4525	4007	3488	2970	584	553	523	492	461	430	369	307	246	184	123	61	0
1500	500	5781	5100	4922	4493	3596	3513	3356	3273	3115	2883	2567	2327	1149	1021	879	426	0
1500	1000	6200	5500	5400	5000	4300	4300	4300	4300	4300	4300	4300	4300	4300	4300	4300	4300	0
1500	1500	7100	6900	6600	6400	6200	6000	5700	5300	5000	4600	4300	4000	3700	3400	3100	2800	0
1500	2000	7400	7100	6700	6500	6300	6000	5700	5300	5000	4600	4300	4000	3700	3400	3100	2800	0
1500	2500	7700	7200	6700	6400	6100	5800	5500	5200	4900	4600	4300	4000	3700	3400	3100	2800	0
1500	3000	8000	7200	6700	6300	5900	5600	5300	5000	4700	4400	4100	3800	3500	3200	2900	2600	0
1500	3500	8300	7200	6700	6250	5700	5400	5100	4800	4500	4200	3900	3600	3300	3000	2700	2400	0
1500	4000	8600	7200	6700	6200	5600	5300	5000	4700	4400	4100	3800	3500	3200	2900	2600	2300	0
1500	4500	8900	7200	6700	6100	5500	5200	4900	4600	4300	4000	3700	3400	3100	2800	2500	2200	0
1500	5000	9200	7200	6700	6000	5400	5100	4800	4500	4200	3900	3600	3300	3000	2700	2400	2100	0
1500	5500	9500	7200	6700	5900	5300	5000	4700	4400	4100	3800	3500	3200	2900	2600	2300	2000	0
1500	6000	9800	7200	6700	5800	5200	4900	4600	4300	4000	3700	3400	3100	2800	2500	2200	1900	0
1500	6500	10100	7200	6700	5700	5100	4800	4500	4200	3900	3600	3300	3000	2700	2400	2100	1800	0
1500	7000	10400	7200	6700	5600	5000	4700	4400	4100	3800	3500	3200	2900	2600	2300	2000	1700	0
1500	7500	10700	7200	6700	5500	4900	4600	4300	4000	3700	3400	3100	2800	2500	2200	1900	1600	0
2000	0	4581	4130	3680	3230	643	609	575	541	507	474	406	338	271	203	135	68	0
2000	500	5870	5420	4970	4708	3636	3752	3594	3510	3352	3118	2801	2710	1193	941	927	434	0
2000	1000	6300	5850	5400	5206	4300	4800	4608	4500	4300	4000	3600	3500	1500	1287	1191	556	0
2000	1500	6700	6300	5800	5700	4600	5100	4900	4800	4600	4300	3900	3800	1600	1400	1311	624	0
2000	2000	7100	6700	6200	6100	5000	5500	5300	5200	5000	4700	4300	4200	1700	1500	1400	779	0
2000	2500	7500	7100	6600	6500	5400	5900	5700	5600	5400	5100	4700	4600	1800	1600	1500	843	0
2000	3000	7900	7500	7000	6900	5800	6300	6100	6000	5800	5500	5100	5000	1900	1700	1600	908	0
2000	3500	8300	7900	7400	7300	6200	6700	6500	6400	6200	5900	5500	5400	2000	1800	1700	973	0
2000	4000	8700	8300	7800	7700	6600	7100	6900	6800	6600	6300	5900	5800	2100	1900	1800	1038	0
2000	4500	9100	8700	8200	8100	7000	7500	7300	7200	7000	6700	6300	6200	2200	2000	1900	1103	0
2000	5000	9500	9100	8600	8500	7400	7900	7700	7600	7400	7100	6700	6600	2300	2100	2000	1168	0
2000	5500	9900	9500	9000	8900	7800	8300	8100	8000	7800	7500	7100	7000	2400	2200	2100	1233	0
2000	6000	10300	9900	9400	9300	8200	8700	8500	8400	8200	7900	7500	7400	2500	2300	2200	1298	0
2000	6500	10700	10300	9800	9700	8600	9100	8900	8800	8600	8300	7900	7800	2600	2400	2300	1363	0
2000	7000	11100	10700	10200	10100	9000	9500	9300	9200	9000	8700	8300	8200	2700	2500	2400	1428	0
2000	7500	11500	11100	10600	10500	9400	9900	9700	9600	9400	9100	8700	8600	2800	2600	2500	1493	0
2000	8000	11900	11500	11000	10900	9800	10300	10100	10000	9800	9500	9100	9000	2900	2700	2600	1558	0
2000	8500	12300	11900	11400	11300	10200	10700	10500	10400	10200	9900	9500	9400	3000	2800	2700	1623	0
2000	9000	12700	12300	11800	11700	10600	11100	10900	10800	10600	10300	9900	9800	3100	2900	2800	1688	0
2000	9500	13100	12700	12200	12100	11000	11500	11300	11200	11000	10700	10300	10200	3200	3000	2900	1753	0
2000	10000	13500	13100	12600	12500	11400	11900	11700	11600	11400	11100	10700	10600	3300	3100	3000	1818	0
2000	10500	13900	13500	13000	12900	11800	12300	12100	12000	11800	11500	11100	11000	3400	3200	3100	1883	0
2000	11000	14300	13900	13400	13300	12200	12700	12500	12400	12200	11900	11500	11400	3500	3300	3200	1948	0
2000	11500	14700	14300	13800	13700	12600	13100	12900	12800	12600	12300	11900	11800	3600	3400	3300	2013	0
2000	12000	15100	14700	14200	14100	13000	13500	13300	13200	13000	12700	12300	12200	3700	3500	3400	2078	0
2000	12500	15500	15100	14600	14500	13400	13900	13700	13600	13400	13100	12700	12600	3800	3600	3500	2143	0
2000	13000	15900	15500	15000	14900	13800	14300	14100	14000	13800	13500	13100	13000	3900	3700	3600	2208	0
2000	13500	16300	15900	15400	15300	14200	14700	14500	14400	14200	13900	13500	13400	4000	3800	3700	2273	0
2000	14000	16700	16300	15800	15700	14600	15100	14900	14800	14600	14300	13900	13800	4100	3900	3800	2338	0
2000	14500	17100	16700	16200	16100	15000	15500	15300	15200	15000	14700	14300	14200	4200	4000	3900	2403	0
2000	15000	17500	17100	16600	16500	15400	15900	15700	15600	15400	15100	14700	14600	4300	4100	4000	2468	0
2000	15500	17900	17500	17000	16900	15800	16300	16100	16000	15800	15500	15100	15000	4400	4200	4100	2533	0
2000	16000	18300	17900	17400	17300	16200	16700	16500	16400	16200	15900	15500	15400	4500	4300	4200	2598	0
2000	16500	18700	18300	17800	17700	16600	17100	16900	16800	16600	16300	15900	15800	4600	4400	4300	2663	0
2000	17000	19100	18700	18200	18100	17000	17500	17300	17200	17000	16700	16300	16200	4700	4500	4400	2728	0
2000	17500	19500	19100	18600	18500	17400	17900	17700	17600	17400	17100	16700	16600	4800	4600	4500	2793	0
2000	18000	19900	19500	19000	18900	17800	18300	18100	18000	17800	17500	17100	17000	4900	4700	4600	2858	0
2000	18500	20300	19900	19400	19300	18200	18700	18500	18400	18200	17900	17500	17400	5000	4800	4700	2923	0
2000	19000	20700	20300	19800	19700	18600	19100	18900	18800	18600	18300	17900	17800	5100	4900	4800	2988	0
2000	19500	21100	20700	20200	20100	19000	19500	19300	19200	19000	18700	18300	18200	5200	5000	4900	3053	0
2000	20000	21500	21100	20600	20500	19400	19900	19700	19600	19400	19100	18700	18600	5300	5100	5000	3118	0
2000	20500	21900	21500	21000	20900	19800	20300	20100	20000	19800	19500	19100	19000	5400	5200	5100	3183	0
2000	21000	22300	21900	21400	21300	20200	20700	20500	20400	20200	19900	19500	19400	5500	5300	5200	3248	0
2000	21500	22700	22300	21800	21700	20600	21100	20900	20800	20600	20300	19900	19800	5600	5400	5300	3313	0
2000	22000	23100	22700	22200	22100	21000	21500	21300	21200	21000	20700	20300	20200	5700	5500	5400	3378	0
2000	22500	23500	23100	22600	22500	21400	21900	21700	21600	21400	21100	20700	20600	5800	5600	5500	3443	0
2000	23000	23900	23500	23000	22900	21800	22300	22100	22000	21800	21500	21100	21000	5900	5700	5600	3508	0
2000	23500	24300	23900	23400	23300	22200	22700	22500	22400	22200	21900	21500	21400	6000	5800	5700	3573	0
2000	24000	24700	24300	23800	23700	22600	23100	22900	22800	22600	22300	21900	21800	6100	5900	5800	3638	0
2000	24500	25100	24700	24200	24100	23000	23500	23300	23200	23000	22700	22300	22200	6200	6000	5900	3703	0
2000	25000	25500	25100	24600	24500	23400	23900	23700	23600	23400	23100	22700	2					

Liquid contact with the heated surface is frequent at low wall superheats but becomes less frequent at higher wall superheats. A similar heat transfer process takes place during low quality convective transition boiling, but here the convective velocities improve the heat transfer coefficient both in film boiling and nucleate boiling.

In the high quality region, most of the heat transferred during transition boiling is due to convection to the vapor and to droplet-wall interaction. Initially, at surface temperatures just in excess of the boiling crisis temperature, a significant fraction of the droplets deposit on the heated surface, but at higher wall superheats, the vapor repulsion forces become significant and repel most of the droplets before they can contact the heated surface.

12B.3.2 Data Trends

A thorough review of the transition boiling literature was conducted by Groeneveld (1976,1977) and Fung (1978). Parametric trends have been deduced from the data (Groeneveld, 1976). In general, an increase in mass flux increases the transition boiling heat flux, especially at higher wall superheats because of its strong positive effects on film boiling heat transfer. The effect of an increase in subcooling is similar but the effect of quality is less clear. The data suggest that at low wall superheats, an increase in quality will have a negative effect on the transition boiling heat flux (similar to its effect on CHF). At higher wall superheats, the film boiling mode dominates and hence, the reverse is true. Fig. 3 illustrates the above trends. Note also the change in slope of the transition boiling curve from a negative value to a positive value at high flows and qualities. This trend is discussed in greater detail by Groeneveld (1980a).

12B.3.3 Correlations

Despite the scarcity of transition boiling data, a large number of correlations have been proposed. They may be divided into three groups:

- (i) Correlations containing boiling and convective components, e.g., Ramu (1974), Mattson (1974) and Tong (1972a). These correlations usually have the form

$$h = A \exp(-B \Delta T_w) + \frac{k_v}{D_e} a \text{Re}_v^b \text{Pr}_v^c$$

where the first term on the right-hand side represents the boiling component (which becomes insignificant at high wall superheats) and second term represents the convective component. These correlations are often claimed to be valid both in the transition boiling and film boiling region.

- (ii) Phenomenological correlations, e.g., Illoeje (1974), Tong and Young (1974). These correlations are based on a physical model of heat transfer in the transition boiling region. Because of an inadequate physical understanding they still contain many empirical constants.

- (iii) Empirical correlations, e.g., Ellion (1954), Berenson (1960) and McDonough (1961). These correlations all have a very simple form and generally cannot be extrapolated outside the range of data on which they are based. Some of these correlations are based on the conditions at the boiling crisis (CHF and T_{CHF}).

12B.3.4 Recommendation

Groeneveld (1976) has shown that frequently (i) the correlations often do not agree with each other, and (ii) their data base is questionable. Hence none of the current correlations can be fully recommended. The most reasonable prediction was obtained by connecting the experimentally determined CHF and minimum film boiling points by a straight line on a log-log plot of q vs. ΔT_w :

$$\frac{q_{TB}}{q_{min}} = \left(\frac{CHF}{q_{min}} \right)^{n_1} \quad \text{where } n_1 \triangleq \frac{\ln(\Delta T_{min}/\Delta T_w)}{\ln(\Delta T_{min}/\Delta T_{CHF})}$$

or

$$\frac{\Delta T_w}{\Delta T_{min}} = \left(\frac{\Delta T_{CHF}}{\Delta T_{min}} \right)^{n_2} \quad \text{where } n_2 \triangleq \frac{\ln(q_{TB}/q_{min})}{\ln(CHF/q_{min})}$$

If experimental values or predictions for q_{min} (to be obtained from $h_{FB} \cdot \Delta T_{min}$), ΔT_{CHF} and ΔT_{min} are questionable or unavailable the following correlation is tentatively suggested:

$$q_{TB} = \max \left(q_{FB}, CHF \left(\frac{\Delta T_{CHF}}{\Delta T_w} \right) \right)$$

where ΔT_{CHF} is obtained from $\Delta T_{CHF} = CHF/h_{CHEN}$ (Chen, 1963), q_{FB} is obtained from an appropriate film boiling model or correlation (Chapter 12B.5), and ΔT_{min} is obtained from the correlation recommended in chapter 12.B.4.

12B.4 Minimum Film Boiling

12B.4.1 General

The minimum film boiling temperature separates the high temperature region where inefficient film boiling or vapour cooling takes place, from the lower temperature region where the much more efficient transition boiling or nucleate boiling occurs. It thus provides a limit to the application of transition boiling and film boiling correlations. In the literature, a large number of terms have been used to describe the phenomenon at the boundary between transition boiling and film boiling, e.g., sputtering, departure from film boiling, rewetting, film boiling collapse, Leidenfrost point, minimum film boiling. Several types of film boiling terminations may be encountered as is shown in Fig. 4.

12B.4.2 Prediction Methods

Groeneveld (1982, 1982a) reviewed the minimum film boiling correlations and divided them into the following groups:

- (i) Equations based on the Taylor-Helmholtz hydrodynamic instability theory, e.g., Berenson (1961), Henry (1974);
- (ii) Equations based on the maximum liquid superheat theory using either an equation-of-state or an equation derived from the classical nucleation analysis, e.g., Spiegler (1963);
- (iii) Empirical correlations based only on minimum flow film boiling data, e.g., Groeneveld (1982).

The theoretically derived equations for $T_{\min}$, such as Berenson's (1960) and Spiegler's (1963), are based on the assumption that the heated surface is isothermal. During the rewetting process, large temperature gradients are set up. If liquid-wall contacts are intermittent, e.g., at the initiation of transition boiling during cooling down of a surface, they will be of very short duration. These contacts are usually not immediately noticeable at the thermocouple junction, which is normally embedded in the heated surface and hence measures an average temperature. The difference between the bulk surface temperature and the true liquid-wall interface temperature can be very significant and has been included in several minimum film boiling correlations. It depends on the ratio $(k \rho C_p)_l / (k \rho C_p)_w$ and can only be ignored for high-conducting surfaces.

12B.4.3 Recommendation

Groeneveld (1982) has compared the prediction trends of various minimum film boiling correlations and noticed significant differences. The large differences in predicted values and parametric trends are not surprising considering the scarcity of the data, and the very different types of film boiling. Despite these differences, the Groeneveld-Stewart (1982) $T_{\min}$ correlation

$$P \leq 9000 \text{ kPa} : T_{\min} = 284.7 + 0.441 P - 3.72 \times 10^{-6} P^2 \\ + \frac{\Delta H \times 10^4}{(2.82 + 0.00122 P) H_{fg}}$$

$$P > 9000 \text{ kPa} : T_{\min} = (\Delta T_{\min, 9000 \text{ kPa}})$$

$$\left(\frac{P_{\text{crit}} - P}{P_{\text{crit}} - 9000} \right) + T_{\text{sat}}$$

is tentatively recommended for the types I-IV film boiling terminations. This recommendation is based on

(i) the approximate agreement of this correlation with the data of Fung (1981), Stewart (1981), Bennett (1966), Bradfield (1967), Lauer (1976), and
(ii) on the correct asymptotic trend of pressure and subcooling. Although in these experiments thin oxide layers were present, no large correction in T_{min} due to oxide layers was needed and, consequently, corrections such as suggested by Henry (1974) are thought to be of secondary importance for flow boiling on engineering surfaces (i.e., Inconel, steel, Zircaloy).

12B.5 Film Boiling

12B.5.1 General

During film boiling the heated surface is cooled by radiation, forced convection to the vapour and by interaction of the liquid and the heated surface. The vapour can become highly superheated; its temperature is controlled both by wall-vapour and vapour-liquid heat exchange. The liquid is thought to be in the form of:

- (a) a dispersed spray of droplets, usually encountered at void fractions in excess of 80%. The corresponding flow regime is often referred to as the liquid-deficient regime because insufficient liquid is available to maintain a wet wall (Fig. 5).
- (b) a continuous liquid core (surrounded by a vapour annulus which may contain entrained droplets) usually encountered at void fractions below 50%. The corresponding flow regime is sometimes referred to as the inverted annular flow regime (Fig. 5).
- (c) a transition between the above two cases, usually in the form of slug flow (Fig. 5).

Of the above post-dryout regimes, the liquid deficient regime is most commonly encountered and has been well studied. Its post-dryout temperature is moderate while for flow regimes (b) and (c) the boiling crisis frequently results in a failure of the heated surface.

12B.5.2 Post-Dryout Models

Post dryout models are models based on the three conservation equations for each phase and suitable empirical interfacial correlations. The first post-dryout models were developed by the UKAEA (Bennett, 1967) and MIT (Laverty, 1967). In these models, all parameters were initially evaluated at the dryout location. It was assumed that heat transfer takes place in two steps: (i) from the heated surface to the vapor, and (ii) from the vapor to the droplets. The models evaluated the axial gradients in droplet diameter, vapor and droplet velocity, and pressure, from the conservation equations. Using a heat balance, the vapor superheat was then evaluated. The wall temperature was finally found from the vapor temperature using a superheated steam heat transfer correlation. Bailey (1972), Groeneveld (1972), and Plummer (1976) have suggested improvements to the original model by including droplet-wall interaction, by permitting a gradual change in average droplet diameter due to the breakup of droplets, and by including vapor flashing for large pressure gradients. Additional expressions for the vapor generation rate have recently been suggested by Saha (1980) and Jones (1977).

Recently models have also been developed for the inverted annular flow regime (Fung, 1981; Kaufman, 1976; Elias, 1981; Chan, 1980). They are basically unequal velocity, unequal temperature (UVUT) models which can

account for the non-equilibrium in both the liquid and the vapor phase.

12B.5.3 Post-dryout Correlations

Empirical post-dryout correlations may be subdivided as follows:

(i) Thermal equilibrium correlations These are correlations which assume that the liquid is in thermal equilibrium with the vapor, and the heated surface is cooled by forced convection to the vapor only. These correlations are basically forced convective correlations where the vapor velocity is evaluated by assuming either homogeneous flow (Dougall, 1963) or by using a suitable slip ratio correlation (e.g., Quinn, 1966).

Thermal equilibrium correlations usually have the form:

$$Nu_v = a \left[\frac{\rho_v \cdot U_v \cdot De}{\mu_v} \right]^b (Pr_v)^c$$

where

$$U_v = \frac{GX}{\rho_v \cdot \alpha} = \frac{G}{\rho_v} \left[X + \frac{\rho_v}{\rho_l} S (1-X) \right]$$

This type of correlation assumes that the liquid and vapor are in equilibrium, i.e., $X = X_e$. This assumption is valid only at high mass flows and high void fractions where the liquid-vapor heat exchange is very efficient.

(ii) Empirical correlations Empirical correlations have been developed for both the liquid deficient regime and the inverted annular flow regime. They generally predict a heat transfer coefficient which is based on the temperature difference between wall and saturation. They are simple to use but have a limited range of validity and should not be extrapolated outside their recommended range.

(iii) Phenomenological correlations These predictive methods are based on the typical mechanisms governing this type of heat transfer and use a heat transfer coefficient based on the vapor temperature. For the liquid deficient regime, phenomenological correlations usually predict the degree of thermal non-equilibrium. Values for the actual vapor temperature T_{va} , or actual quality, X_a , can be generated from existing post-dryout data, provided the assumption of forced convective cooling only (i.e., no droplet-wall interaction) of the heated wall in the post-dryout region is correct. This approach has been used by Tong (1974), Plummer (1976), Groeneveld-Delorme (1976a), Chen (1979) and Shah (1980).

Groeneveld (1983) has presented a summary of recent non-equilibrium correlations. In general, these correlations are a significant improvement over the empirical post-CHF correlations. They can be given the correct asymptotic trends to let them smoothly converge with the single-phase superheated-steam cooling correlations.

Attempts have also been made to develop phenomenological correlations for the inverted annular regime (Dougall, 1963; Dalinin, 1969, 1970) and the slug flow film boiling regime (Chi, 1967; Kalinin, 1970). Due to a limited data base, they have not yet been verified.

(iv) Pool film boiling correlations These correlations are applicable in pool boiling and low mass velocities. Most of these correlations basically have the form:

$$h_c = A \left[\frac{g k_v^3 \rho_v (\rho_l - \rho_v) H'_{fg}}{\mu_v \Delta T_s} \right]^{\frac{1}{4}} f(U, \lambda) + h_{rad}$$

which was originally derived by Bromley (1950) from Nusselt's classical analysis of filmwise condensation. The latent heat H'_{fg} is usually modified to include the vapor superheat, while the velocity effect is included theoretically or empirically through $f(U, \lambda)$ where λ is either equal to the diameter, the critical wavelength or the most unstable wavelength.

(v) No evaporation after dryout In addition to the above prediction methods, one can also assume that, in the liquid deficient regime, no evaporation will take place, i.e., the wall heat flux is used only for superheating of the vapor and X_a remains constant. Here the predicted wall temperature is very high since the vapor becomes progressively more superheated. This prediction is pessimistic except at lower flows (Bennett, 1967). It is, however, useful as an upper boundary for the heated surface temperature.

12B.5.4 Discussion

During the past fifteen years, much progress has been made in understanding the physical mechanisms governing post-dryout heat transfer. Post-dryout models have been reasonably successful in predicting the post-dryout temperature distribution. These models, however, require accurate values of the dryout quality; they are generally considered useful especially for predictions outside the data base (provided the assumed physical mechanisms are still valid).

Of the prediction methods discussed in the previous section, the phenomenological correlations are considered the most accurate. They tend to satisfy the observed experimental data trends, hence they have a much wider range of applicability than the empirical post-dryout correlations, but they are generally more difficult to use. Figure 6 shows schematically the relative temperature predictions of four of the prediction methods. Note that curve A, based on the "thermal equilibrium after dryout" assumption, represents a lower boundary while curve B, based on the "no evaporation after dryout" assumption represents an upper boundary for the post-dryout temperature.

12B.6 References

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Nomenclature

Cp	Specific heat at constant pressure	$\text{kJ} \cdot \text{kg}^{-1} \cdot \text{K}^{-1}$
CHF	Critical heat flux	$\text{kW} \cdot \text{m}^{-2}$
D	Tube diameter	m
De	Hydraulic equivalent diameter	m
g	Acceleration due to gravity	$\text{m} \cdot \text{s}^{-2}$
G	Mass flux	$\text{kg} \cdot \text{m}^{-2} \cdot \text{s}^{-1}$
H	Enthalpy	$\text{kJ} \cdot \text{kg}^{-1}$
h	Heat transfer coefficient	$\text{kW} \cdot \text{m}^{-2} \cdot \text{K}^{-1}$
k	Thermal conductivity	$\text{kW} \cdot \text{m}^{-1} \cdot \text{K}^{-1}$
Lh	Heated length	m
Nu	Nusselt number ($= h \cdot D / k$)	-
P	Pressure	kPa
Pr	Prandtl number ($= \mu \cdot \text{Cp} / k$)	-
q	Surface heat flux	$\text{kW} \cdot \text{m}^{-2}$
Re	Reynolds number ($= \rho \cdot U \cdot D / \mu$)	-
S	Slip ratio ($= U_v / U_\ell$)	-
T	Temperature	$^{\circ}\text{C}$
U	Velocity	$\text{m} \cdot \text{s}^{-1}$
X	Flow quality (vapor weight fraction)	-

Greek

α	Void fraction	-
ΔH	Local subcooling	$\text{kJ} \cdot \text{kg}^{-1}$
ΔT	Temperature difference (usually with respect to saturation)	K
μ	Viscosity	$\text{kg} \cdot \text{m}^{-1} \cdot \text{s}^{-1}$
ρ	Density	$\text{kg} \cdot \text{m}^{-3}$
σ	Surface tension	$\text{N} \cdot \text{m}^{-1}$

Subscripts

a	Actual
c	Critical
do	Value pertaining to onset of dryout condition
e	Equilibrium
f	Saturated liquid
fg	Difference between saturated liquid and saturated vapor
g	Saturated vapor
in	Inlet
l, l	Liquid (may be subcooled)
max	Maximum
min, mfb	Minimum film boiling value
Q	Apparent quench temperature
s, sat	Saturated
sub	Subcooling
v	Vapor
w	Heat wall

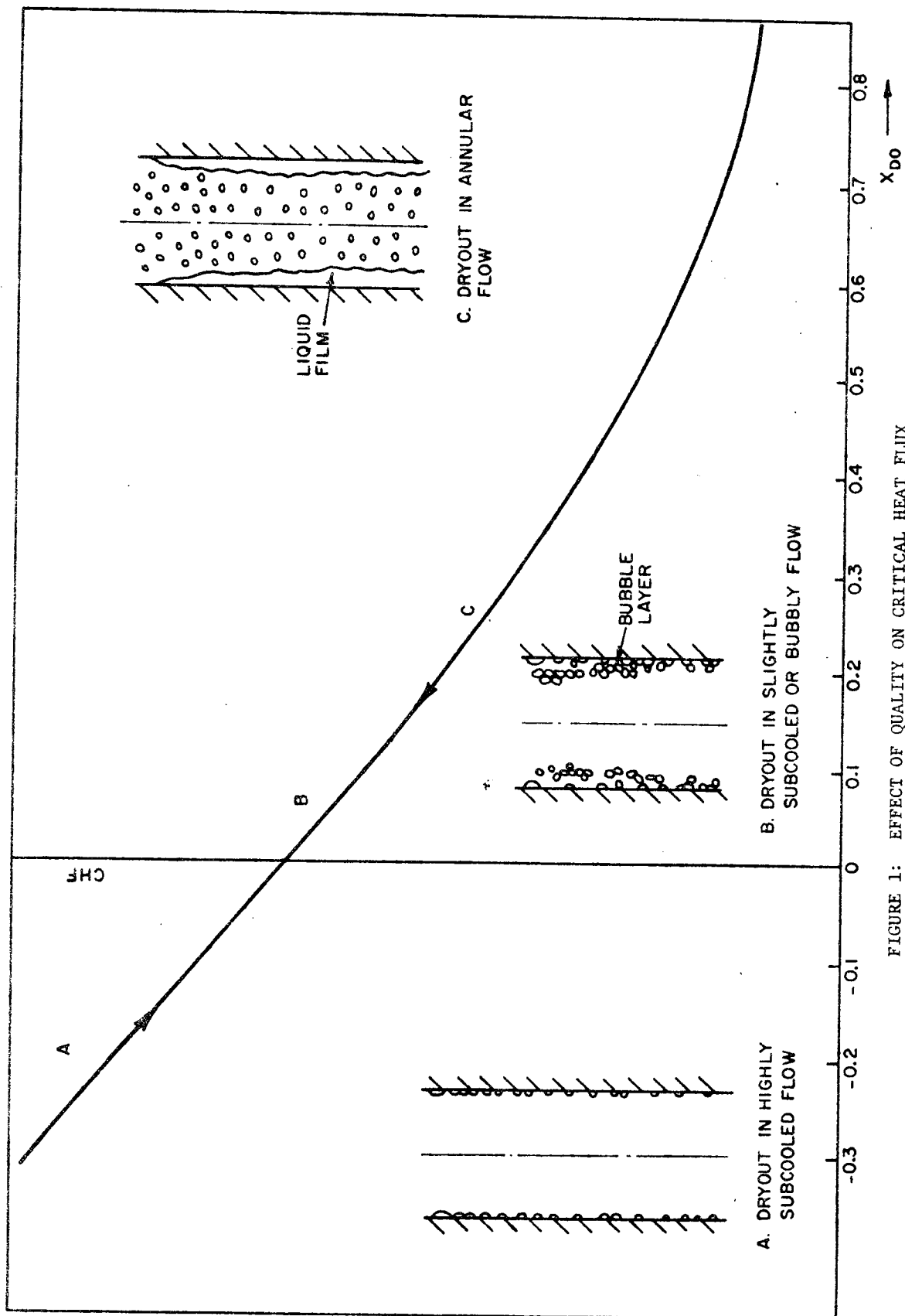


FIGURE 1: EFFECT OF QUALITY ON CRITICAL HEAT FLUX

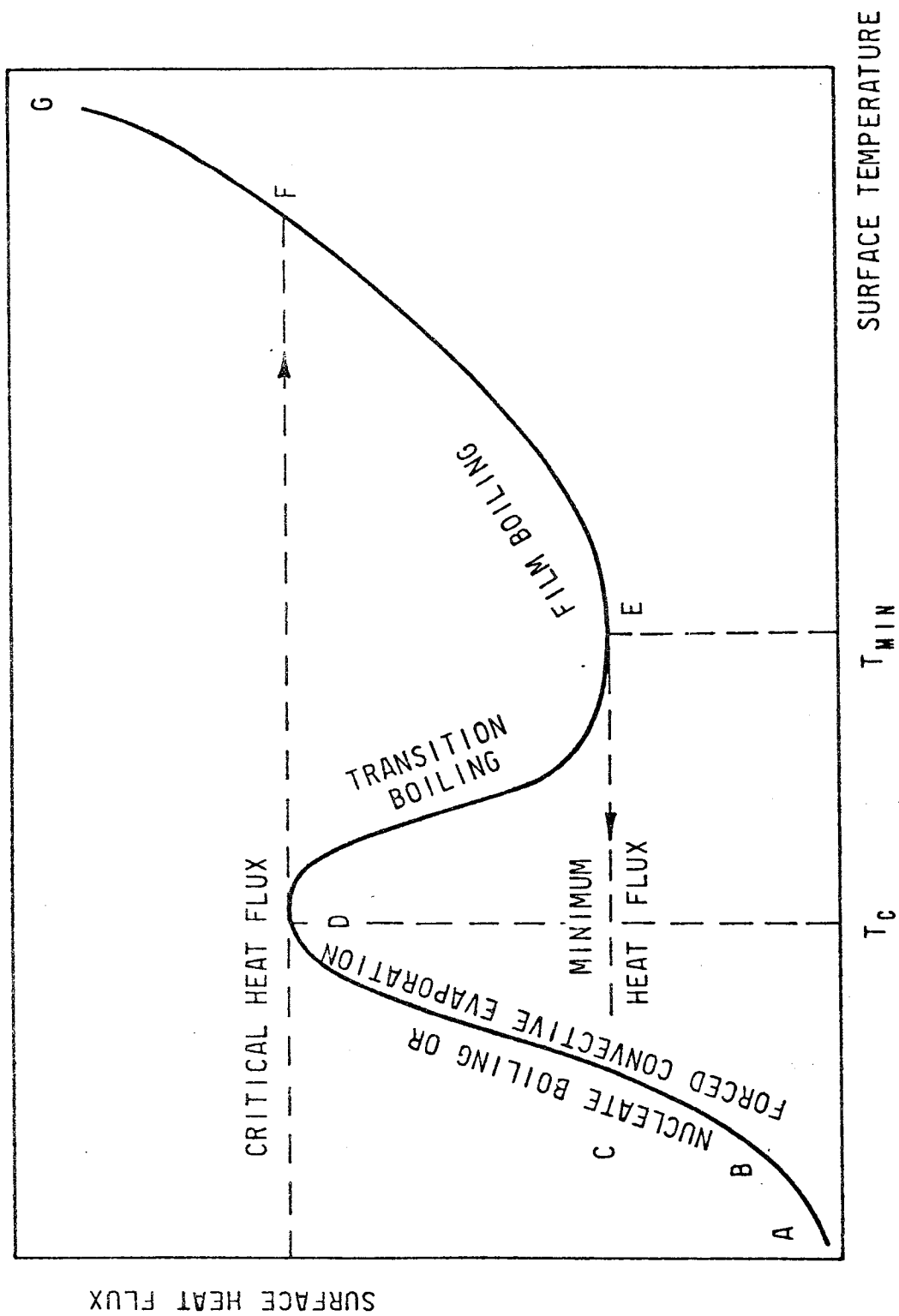
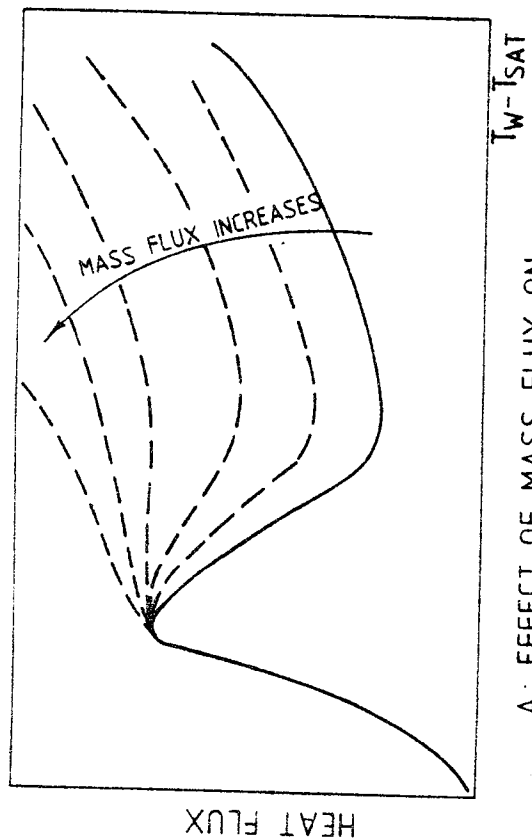
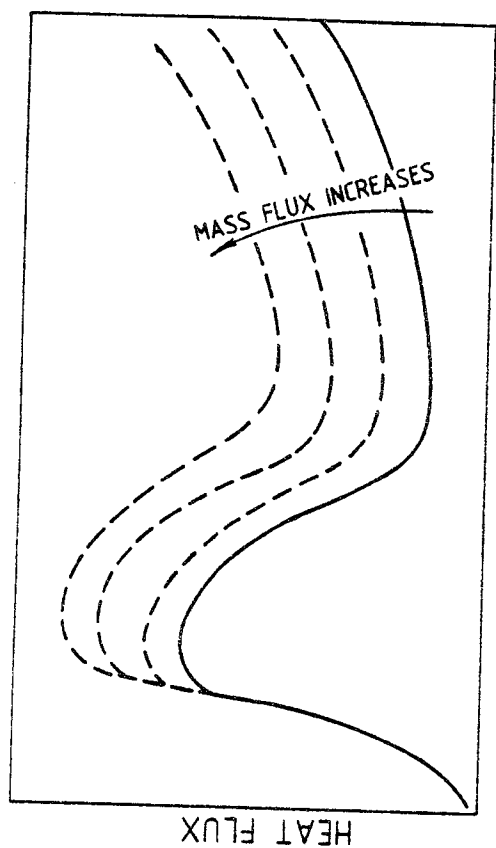


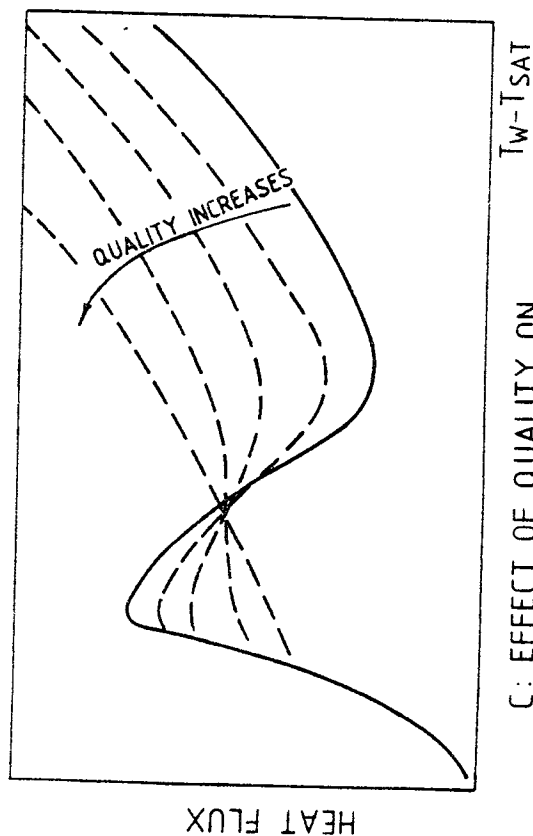
FIG.2 CONVECTIVE BOILING CURVE



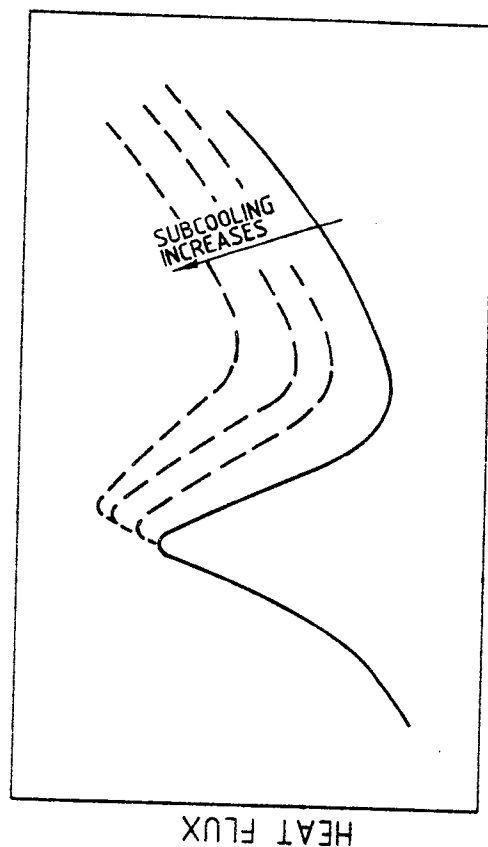
A: EFFECT OF MASS FLUX ON
BOILING CURVE (ANNULAR FLOW REGIME)



B: EFFECT OF MASS FLUX ON
BOILING CURVE (FORCED CONVECTIVE
SUBCOOLED OR LOW QUALITY BOILING)

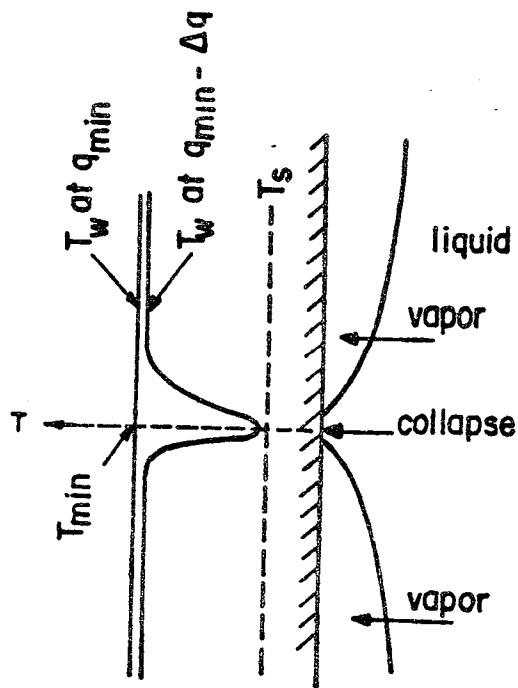


C: EFFECT OF QUALITY ON
FORCED CONVECTIVE BOILING CURVE

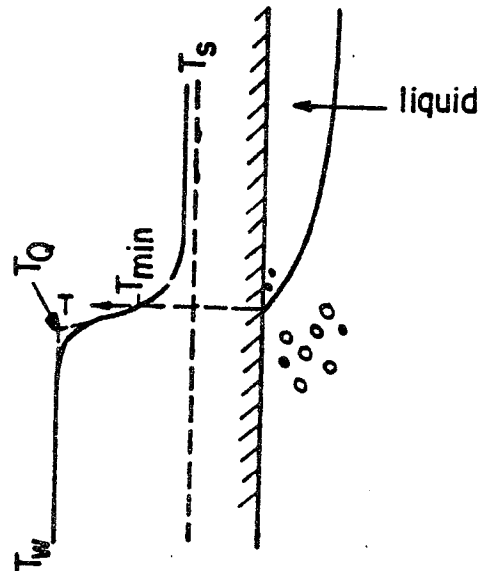


D: EFFECT OF SUBCOOLING ON
BOILING CURVE

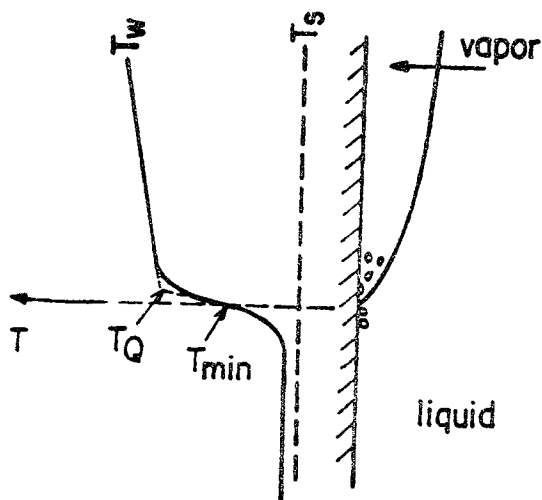
FIG. 3 PARAMETRIC TRENDS OF FLOW BOILING CURVE



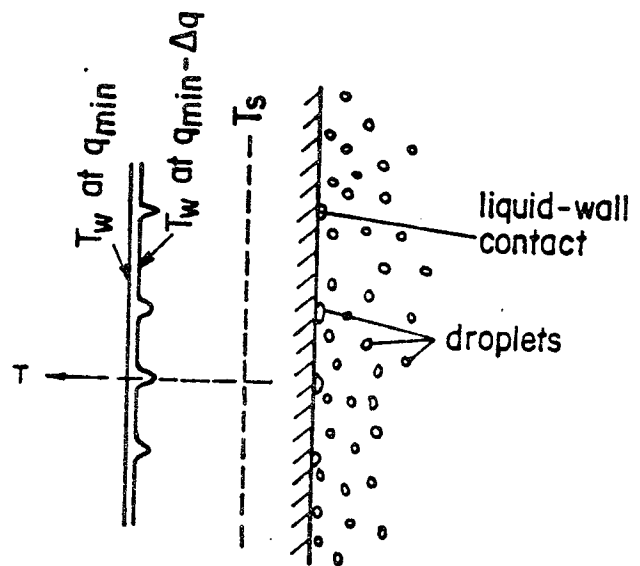
Type I : Collapse of vapor film



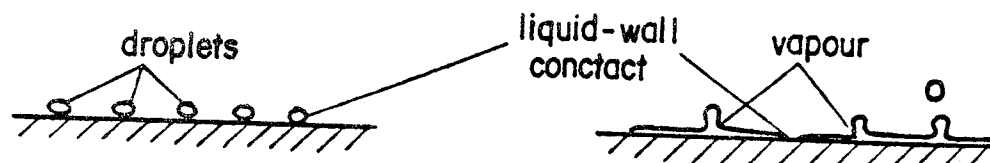
Type II : Top flooding



Type III : Bottom flooding



Type IV : Droplet cooling



Type V : Leidenfrost boiling

Type VI : Pool boiling

FIG. 4: FILM BOILING TERMINATION TYPES

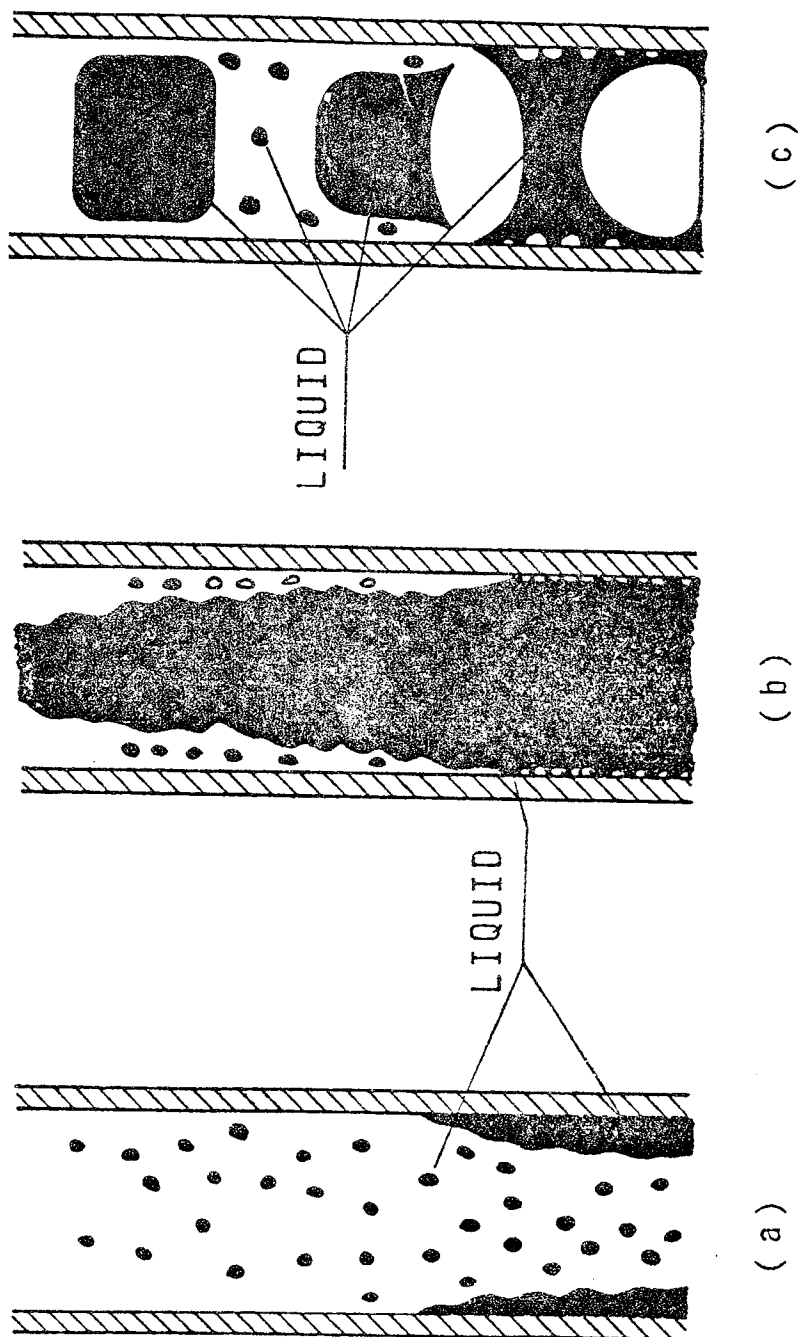
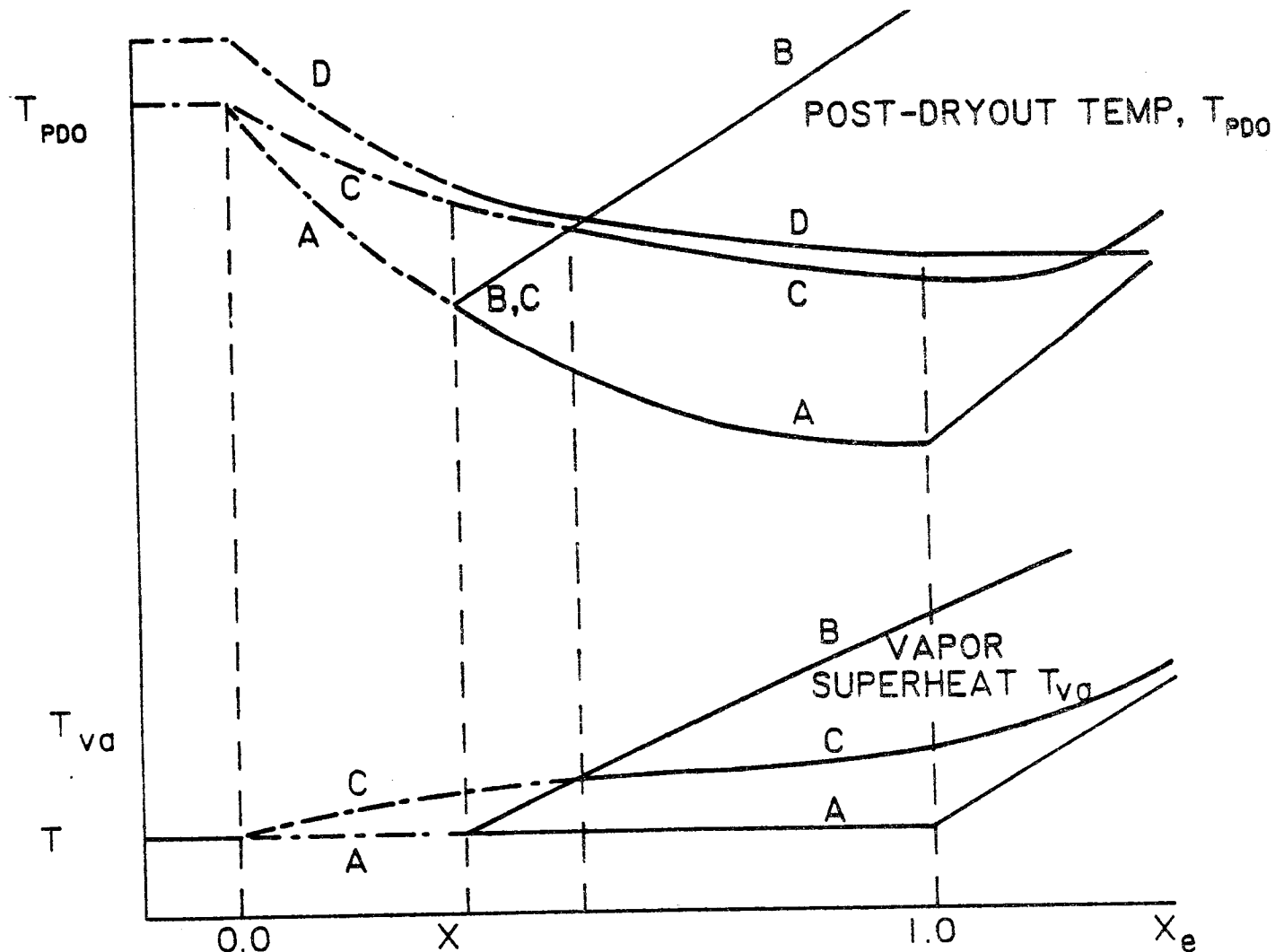


FIG.5: FLOW REGIMES DURING FILMBOILING
 (a) LIQUID DEFICIENT FLOW
 (b) INVERTED ANNULAR FLOW
 (c) SLUG FLOW



- A - THERMAL EQUILIBRIUM
- B - NO EVAPORATION AFTER DRYOUT
- C - THERMAL EQUILIBRIUM + CORRELATION FACTOR
- D - EMPIRICAL POST-DRYOUT CORRELATION EXTENSION OF PREDICTION IF X_{do} IS UNKNOWN

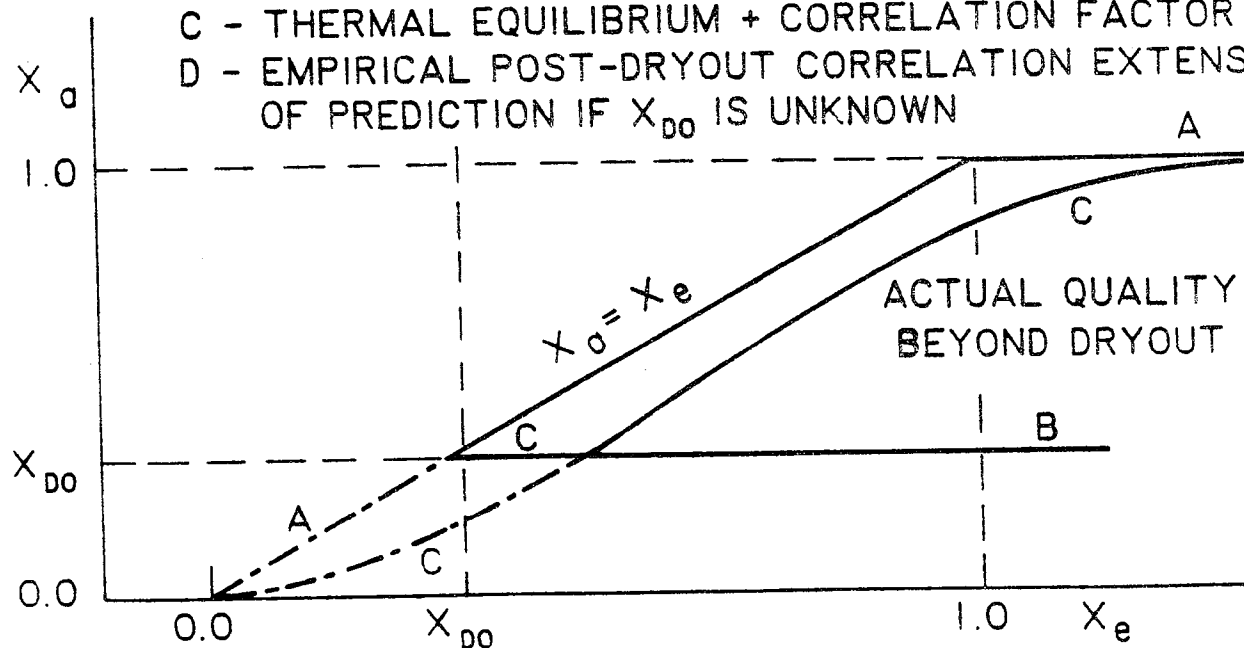


FIGURE 6 : VARIOUS PREDICTIONS OF POST-DRYOUT PARAMETERS

